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Muon spin rotation and neutron scattering investigations of the *B*-site ordered double perovskite Sr₂DyRuO₆

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The magnetic ground state of the *B*-site ordered double perovskite Sr₂DyRuO₆ has been investigated using muon spin rotation and relaxation (μ SR), neutron powder diffraction (NPD), and inelastic neutron scattering (INS), in addition to heat capacity and magnetic susceptibility (AC and DC) measurements. A clear signature of a long-range ordered magnetic ground state has been observed in the heat capacity data, which exhibits two sharp anomalies at 39.5 and 36 K found as well in the magnetic data. Further supporting evidence consistent with long-range magnetic ordering comes from a sharp drop in the muon initial asymmetry and a peak in the relaxation rate at 40 K, along with a weak anomaly near 36 K. Based on temperature dependent NPD, the low temperature magnetic structure contains two interpenetrating lattices of Dy³⁺ and Ru⁵⁺, forming an antiferromagnetic ground state below 39.5 K with magnetic propagation vector $k = (0, 0, 0)$. The magnetic moments of Dy³⁺ and Ru⁵⁺ at 3.5 K are pointing along the crystallographic *b* axis with values of $\mu^{\text{Dy}} = 4.92(10) \mu_B$ and $\mu^{\text{Ru}} = 1.94(7) \mu_B$, respectively. The temperature dependence of the Ru⁵⁺ moments follows a mean field type behavior, while that of the Dy³⁺ moments exhibits a deviation indicating that the primary magnetic ordering is induced by the order of the 4*d* electrons of Ru⁵⁺ rather than that of the proper 4*f* Dy³⁺ electrons. The origin of the second anomaly observed in the heat capacity data at 36 K must be connected to a very small spin reorientation as the NPD studies do not reveal any clear change in the observed magnetic Bragg peaks' positions or intensities between these two transitions. INS measurements reveal the presence of crystal field excitations (CEF) in the paramagnetic state with overall CEF splitting of 73.8 meV, in agreement with the point charge model calculations, and spin wave excitations below 9 meV at 7 K. Above T_N , the spin wave excitations transform into a broad diffuse scattering indicating the presence of short-range dynamic magnetic correlations.

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I. INTRODUCTION

Geometrically frustrated antiferromagnetic (AFM) materials have attracted considerable interest over the past few years, motivated by their tendency to form rather exotic magnetic ground states such as the spin glass, spin liquid, or spin ice instead of long-range magnetic order in apparent defiance of the third law of thermodynamics [1–6]. Among the four “canonical” geometrically frustrated lattices: triangular planar, kagome, pyrochlore, and face-centered cubic (fcc), the latter has recently gained strong attention [6–9]. In real materials, the fcc magnetic lattice is conveniently realized in the *B*-site ordered double perovskites, A₂BB'O₆ [10]. Here a magnetic ion resides on the *B'* site, while *B* can be either

magnetic or nonmagnetic and *A* site is nonmagnetic. Both the *B* and *B'* sites constitute interpenetrating face-centered cubic sublattices in which, if the exchange constraints between nearest neighbors are AFM, the basic criteria for geometric frustration are satisfied [6,9,11–15].

Recently, the double perovskites compounds with general formula A₂BB'O₆, with *A* alkaline metals, *B* rare earth metals, and *B'* transition metal, have attracted considerable attention due to their interesting physical properties as well as possible applications in renewable energy and spintronic devices [6,10–17]. Within this class of materials, there are compounds with properties such as a high Curie temperature T_C [18,19], phase separation [20], a high magnetoresistance [21,22], a metal–insulator transition [23,24], and half-metallic antiferromagnets [16,21]. Besides the interesting fundamental physics, double perovskite materials are important for optoelectronic applications and technology [25]. This huge variety of properties has its origin in the possibility of doping and substituting the perovskite structure at the *A* and *B* sites,

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allowing tailoring of the electronic, crystal, and magnetic structure of the compounds, which, in turn, interact with each other. $\text{Sr}_2\text{FeMoO}_6$ was the first double perovskite for which a high magnetoresistance at room temperature was reported ($T_C \sim 420$ K) [26]. By electron doping in similar compounds, the Curie temperatures rises to 635 K for $\text{Sr}_2\text{CrReO}_6$ [27–29] and 750 K for $\text{Sr}_2\text{CrOsO}_6$ [30] which is so far the highest T_C observed in ferrimagnetic double perovskites. A special type of double exchange interactions [31,32] was shown to be responsible for the high magnetic transition temperatures and the strong spin polarization in double perovskites where B and B' cations are in a mixed valence state [33]. Adoption of integer valences leads to reduced T_C or to antiferromagnetic order [33,34].

Among the antiferromagnetically ordered double perovskites, $\text{Sr}_2\text{LnRuO}_6$ (Ln = rare earth, Y, Ho, Yb, Lu, etc.) compounds exhibit many interesting properties, for example the presence of two magnetic phase transitions and strong geometrical frustration above the magnetic ordering up to as high as 300 K, confirmed via heat capacity and inelastic neutron scattering measurements, respectively [35–39]. Recent neutron diffraction studies on Sr_2YRuO_6 reveal that at the first transition temperature only half of the Ru layers order magnetically while the other half (alternatively) reveals short-range ordering and below the second phase transition the system exhibits a type-I AFM ground state [36]. Although the presence of frustration has been observed in many double perovskite compounds, its origin is not clear at present. In addition, diffuse scattering has been observed in the compounds Sr_2YRuO_6 , $\text{La}_2\text{NaRuO}_6$, $\text{La}_2\text{NaOsO}_6$, and $\text{Sr}_2\text{YbRuO}_6$ [36,37,40] in which $\text{La}_2\text{NaRuO}_6$ reveals a single magnetic transition below 15 K to an incommensurate magnetic ground state, whereas $\text{La}_2\text{NaOsO}_6$ does not exhibit any long-range order down to 4 K on the quasi-fcc lattices as a result of geometrical frustration [11,40]. These results motivate the investigation of other double perovskite compounds in order to understand the presence of geometrical frustration and its effect on the magnetic ground state. We have therefore studied the detailed dynamic and static magnetic properties of $\text{Sr}_2\text{DyRuO}_6$ (SDRO) using magnetization, heat capacity, muon spin resonance/rotation (μSR), neutron powder diffraction (NPD), and inelastic neutron scattering (INS) measurements. SDRO exhibits a magnetic anomaly at ~ 40 K, which is suspected to be associated with the long-range ordering [41,42]. An exchange bias effect in SDRO has also been observed below the AFM ordering temperature and the possible cause for the observed effect was suspected to be linked to Dzyaloshinskii-Moriya (DM) interactions present in this geometrically frustrated system [41]. DFT results report that the main contribution to the spin moment comes from the f orbitals, with a considerable role of the d orbitals, and suggest that SDRO will behave as a conductor and semiconductor for spin-up and spin-down orientations, respectively [43]. However, no direct evidence or studies about the electronic or magnetic structure/ground state is available on SDRO in the existing literature. The present work will fill the gap to understand the low temperature magnetic behavior of SDRO and provide an ideal example to compare with the available data of other geometrically frustrated double perovskites having two magnetic cations at the B and B' sites.

II. EXPERIMENTAL DETAILS

The polycrystalline sample (10 g) of $\text{Sr}_2\text{DyRuO}_6$ (SDRO) was prepared by solid-state reaction from stoichiometric amounts of SrCO_3 , RuO_2 , and Dy_2O_3 (Aldrich 99.99%) which were mixed in an agate mortar and pestle, and pressed into pellets. These pellets were then annealed at 1123 K for 12 h and sintered at 1253 K for 24 h, with frequent regrinding and repelletizing. The structure characterization at 300 K was carried out using the GEM time-of-flight (TOF) neutron powder diffractometer (NPDM) at the ISIS neutron Facility, UK. The DC-magnetic susceptibility and magnetization isotherm were measured using a SQUID magnetometer (Quantum Design, MPMS). Heat capacity measurements were performed using a relaxation technique by a commercial system (Quantum Design, PPMS) in the temperature range of 1.8–100 K. The AC susceptibility was measured using the same Quantum Design PPMS. To investigate the magnetic structure/ground state, low temperature NPD measurements were performed using the constant wavelength ($\lambda = 2.396$ Å) high intensity diffractometer D20 between 1.7 and 50 K at ILL Grenoble, France. High-resolution data were recorded as well at the ILL on the powder diffractometer D2B using $\lambda = 1.594$ and 2.399 Å. All the diffraction data have been analyzed using the Rietveld refinement program Fullprof [44]. The μSR experiments were carried out using the MuSR spectrometer in the longitudinal geometry at the ISIS muon facility, UK. We have performed zero-field (ZF) and longitudinal-field (LF) μSR measurements between 1.5 and 50 K and LF field between 0 and 2500 G. The powder sample (thickness ~ 2 mm) was mounted onto a 99.995+ % pure silver plate using GE varnish and was covered with 18 μm silver foil. Inelastic neutron scattering measurements were performed on the time-of-flight spectrometers MERLIN at ISIS Facility and IN6 at ILL, Grenoble. We used a powder sample of SDRO in an annular Al can of outer diameter 40 mm on MERLIN and 20 mm on IN6.

III. RESULTS AND DISCUSSION

A. Room temperature structural characterization

Figure 1 shows the NPD pattern of SDRO collected at 300 K from the 34.96° detector bank of the GEM diffractometer. The structure was refined using the monoclinic space group $P2_1/n$, assuming a 1:1 ordering of the Dy^{3+} and Ru^{5+} cations. The Dy^{3+} and Ru^{5+} cations occupy distinct Wyckoff sites $2c$ and $2d$, respectively, resulting in the ordered arrangement. No impurity peaks were detected within the resolution limit. The refined lattice parameters at 300 K are $a = 5.7774(2)$ Å, $b = 5.7948(2)$ Å, $c = 8.1848(2)$ Å, $\beta = 90.181(3)^\circ$, and $V = 276.88(1)$ Å³. The determined lattice and structural parameters are in good agreement with the existing literature [41–43]. The refinement did not give any evidence for a possible site disorder between the Dy^{3+} and Ru^{5+} cations. Therefore, our results confirm the ordered double perovskite structure of SDRO.

B. Physical properties

Figure 2 shows the measured heat capacity of SDRO as a function of temperature for zero field (ZF) and in applied

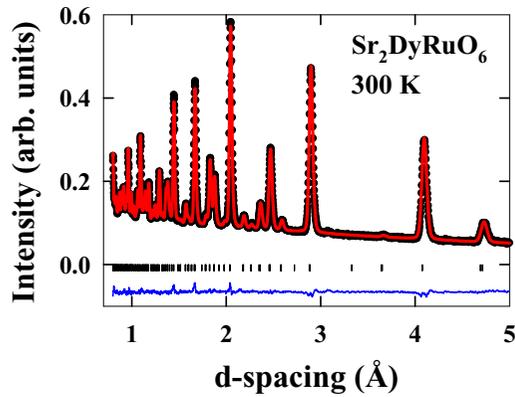


FIG. 1. Rietveld fit to NPD patterns collected at 300 K on GEM diffractometer at ISIS: Black crosses show observed data points; the red line shows the calculated profile; the lower blue line is the difference profile (obs. – calc.); black vertical markers indicate Bragg peak positions.

fields of 1 and 9 T. Two anomalies are evident at 39.5 and 36 K (see the inset in Fig. 2) in the ZF heat capacity data, which disappear in a field of 9 T. The anomalies are more clearly visible in the first order derivative which is presented as an inset in the same figure. At 1 T field, the lower transition does not change much, but the higher transition broadens and moves towards higher temperature. Similar two anomalies/transitions were also reported for isostructural $\text{Sr}_2\text{LnRuO}_6$ ($\text{Ln} = \text{Y}, \text{Ho}, \text{Yb}, \text{and Tb}$) and identified as antiferromagnetic ordering temperatures (T_{N2} and T_{N1}). The anomalies were situated at 24 and 29 K for Y [35,36], 36 and 40 K for Yb [37], 32 and 26 K for Lu [45], and 15 and 36 K for Ho-based systems [38,46]. On the other hand, the heat capacity study on the cubic $\text{Ba}_2\text{DyRuO}_6$ reveals only a single anomaly at 47 K [47], similar to $\text{La}_2\text{NaRuO}_6$ [11,40]. Furthermore, $\text{Sr}_2\text{FeOsO}_6$ exhibits two magnetic transitions at $T_{N1} = 140$ K and $T_{N2} = 67$ K, where both the Fe and Os moments order and the second transition is associated with the change in magnetic structure from AF1 to AF2 [48]. Considering the observation of a spin

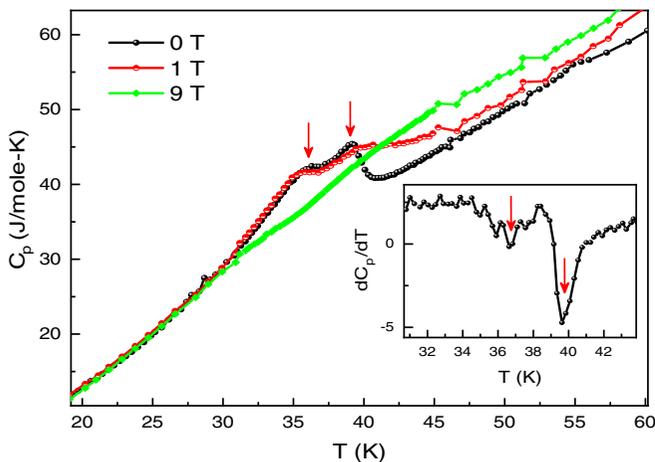


FIG. 2. Heat capacity measurements in the low temperature range measured in zero field and applied field of 1 and 9 T. The inset shows dC_p/dT vs T near the magnetic ordering.

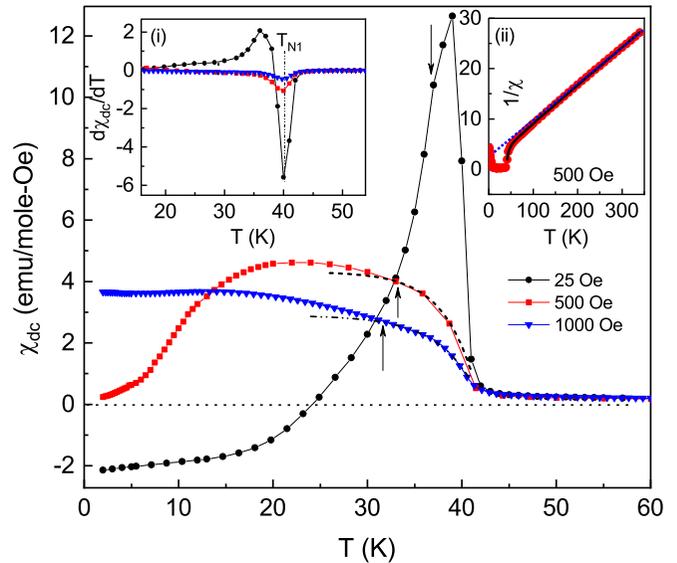


FIG. 3. (a) DC magnetic susceptibility (χ_{DC}) measured at various applied magnetic fields. The dashed and dotted lines are guides to the eye. Inset (i) shows the first derivative of ZFC χ_{DC} and the inset (ii) shows the inverse susceptibility ($1/\chi_{\text{DC}}$) measured in an applied field of 500 Oe. The dotted blue line shows the fit to Curie-Weiss behavior and the solid black line shows the fit using the two-level paramagnetic model (see text). It is to be noted that the negative magnetization seen below 25 K in the 25 Oe data is an artifact due to trapped field in the superconducting magnet of the SQUID magnetometer [74].

gap only below T_{N2} in the inelastic study of $\text{Sr}_2\text{FeOsO}_6$, it was suggested that spin-orbit coupling is important for ground state selection in this compound [48]. This suggests that the two anomalies observed in the heat capacity of SDRO are possibly associated either to the separate long-range magnetic ordering of the Ru and Dy moments or to a spin reorientation transition.

The temperature dependence of the DC-magnetic susceptibilities (χ_{DC}) of SDRO in various applied magnetic fields is shown in Fig. 3 in zero-field cooled (ZFC). The increase of susceptibility below 42 K, irrespective of the applied field value, indicates the emergence of long-range magnetic ordering. With further decreasing temperature, χ_{DC} first increases and exhibits a sharp peak near 40 K for $B = 25$ Oe, nearly matching the first anomaly observed in the heat capacity data (39.5 K). For $B = 500$ Oe, the peak in the susceptibility becomes quite broad exhibiting a plateau. For $B = 1000$ Oe, there is no visible peak and the susceptibility keeps on increasing down to 2 K. As it is difficult to identify the magnetic ordering temperature directly from the χ_{DC} behavior, the first derivative of χ_{DC} is plotted in the left inset of the same figure and shows for all three field values a clear peak at 40 K. This is in accordance with the heat capacity results where the first anomaly was observed at $T_{N1} = 39.5$ K. No direct signature of a second anomaly as found in the heat capacity data ($T_{N2} = 36.5$ K) is evident from the susceptibility data (Fig. 3) for $B = 500$ and 1000 Oe while the derivative points to $T_{N2} = 36$ K for $B = 25$ Oe (inset of Fig. 3). Only an indirect indication of T_{N2} can be found by the rate of change of χ_{DC} for

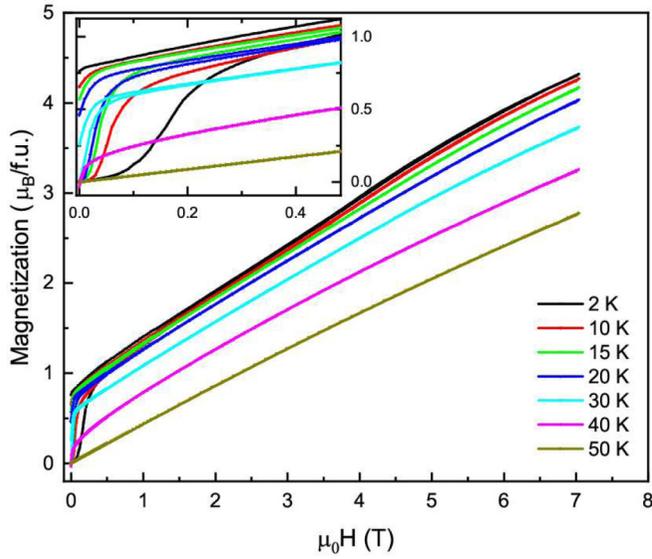


FIG. 4. Magnetization isotherms measured at various temperatures ranging from 2 to 50 K. The inset shows the enlarged view at lower fields. A clear hysteresis can be seen for $T < 40$ K.

the $B = 500$ and 1000 Oe field curves which changes below 35 K. In Fig. 3 the changes in slopes of the magnetic susceptibility are presented by the dashed black lines and using black arrows. As a function of the applied field strength, the values of T_{N2} can be estimated as 36.7, 34, and 32 K for $B = 25, 500$, and 1000 Oe, respectively. The Curie-Weiss (CW) fit of the inverse susceptibility for 500 Oe data is also shown in the same figure in the right inset (dotted blue line), which gives $C = 13.56$ (K emu/mol Oe) and a CW-paramagnetic temperature of $\theta_{CW} = -27.5$ K. The estimated total value of the effective paramagnetic moment is $10.42 \mu_B$ which is slightly smaller than the theoretical value arising from the combined paramagnetic contribution of Dy^{3+} and Ru^{5+} ions which amounts to $11.33 \mu_B$ ($\mu_{eff} = \sqrt{(\mu_{eff}^{Dy^{3+}})^2 + (\mu_{eff}^{Ru^{5+}})^2}$). Furthermore, we have also fitted the inverse susceptibility of SDRO using a two-sublattice paramagnetic model [49,50], $\chi^{-1}(T) = T/C_m + 1/\chi_0 - b/(T - \theta_p)$, which gives an excellent agreement with the data above T_{N1} (see solid line in inset (ii) of Fig. 3), with fitting parameters $C_m = 13.30$ K emu/mol, $\chi_0 = 0.58$ (emu/mol), $b = 10.59$ K mol/emu, and $\theta_p = 37.88$ K.

The isothermal magnetization behavior (M vs H) of SDRO is presented in Fig. 4 as a function of applied magnetic field (H) at selected temperatures ranging from 2 to 50 K. To perform these measurements, the sample was cooled each time from the paramagnetic state (300 K) to avoid any magnetic history effect. The M vs H isotherm at 50 K is almost linear in H , as expected for a paramagnet state. At 2 K, the M vs H curve initially increases rapidly with increasing field up to 0.4 T before it exhibits an almost linear field dependence. The rapid increase in magnetization is also observed for temperatures between 2 and 40 K, but the value above which the field dependence is showing a linear behavior is reducing with increasing temperature. The observed weak ferromagnetic-type behavior in M vs H data at low fields in the antiferromagnetic state, which is also observed in other Sr_2LnRuO_6 ($R =$ rare earth) compounds [37–39], has been attributed to

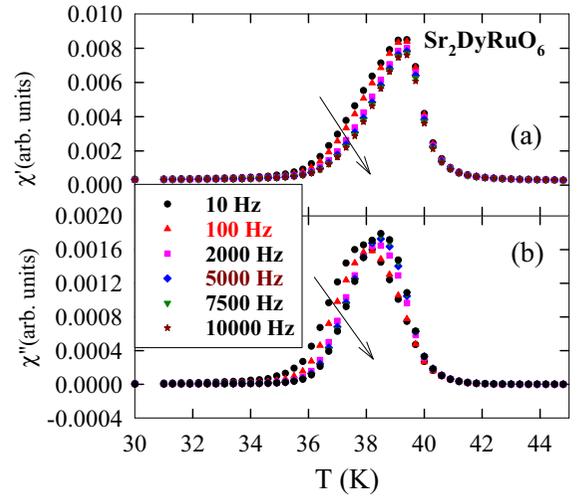


FIG. 5. (a) Real and (b) imaginary components of AC susceptibility (χ_{AC}) measured at various frequencies as a function of temperature.

the contribution of the weak ferromagnetic component from the DM interaction. The magnetization attains a value of only $\sim 4.31 \mu_B$ at 7 T and it does not exhibit saturation behavior at 7 T field at 2 K, which is in good agreement with existing literature [41].

Figure 5 represents the real χ'_{AC} and the imaginary part χ''_{AC} of the temperature dependent AC susceptibility of SDRO at frequencies ranging from 10 Hz to 10 kHz. Both anomalies as observed in the heat capacity and χ_{DC} are visible in the real and imaginary parts. The anomaly at $T_{N1} = 40$ K is revealed by a sharp jump in both χ'_{AC} and χ''_{AC} , and is frequency independent, whereas the anomaly at T_{N2} creates a broad shoulder to the main transition and shows a weak frequency dispersion. This indicates that the second ordering is possibly associated with a very small change in the spin structure near T_{N2} .

C. μ SR measurements

In order to gain insight as to whether the two observed phase transitions in the heat capacity originate from the magnetic ordering, we have investigated SDRO using the μ SR technique. μ SR is a local microscopic probe and is sensitive to extremely small internal fields and is ideal to detect spatially inhomogeneous magnetic features. It is extensively applied to investigate small changes in magnetism [51]. It is interesting to mention that μ SR studies on double perovskites have provided important information on the magnetic ground state of these systems [40,52–54], including information on the microscopic coexistence of magnetic and nonmagnetic phases in $Ba_2PrRu_{0.9}Ir_{0.1}O_6$ [55]. For the present study, zero field (ZF) and longitudinal field (at constant temperature) μ SR measurements have been performed. Figure 6 shows the muon asymmetry versus time spectra at several temperatures between 5 and 90 K measured in ZF. The analysis of μ SR spectra was carried out using a stretched exponential function with constant background:

$$G_z(t) = A_0 \exp[-(\lambda t)^\beta] + A_{bg}. \quad (1)$$

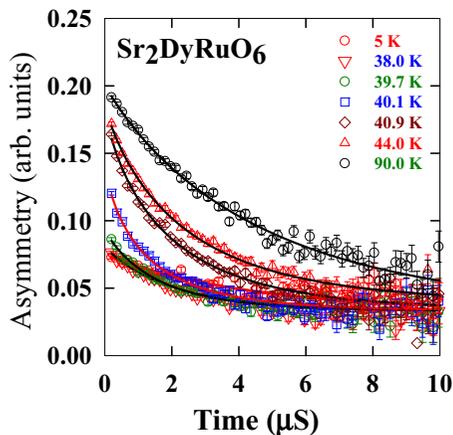


FIG. 6. Zero-field μ SR spectra measured at various temperatures. The experimental data are shown by the symbols and the solid line shows fit to the data using stretch exponent function.

Here A_0 is the muon initial asymmetry and λ is the muon relaxation rate. If the exponent $\beta = 2$, the function becomes Gaussian while for $\beta = 1$, the function becomes exponential. A_{bg} is the constant background arising from muons stopping on the Ag-sample holder. Furthermore, $\beta < 1$ describes inhomogeneous dynamic relaxation where the relaxation is locally exponential but the local rates are distributed [56]. It has no basic theoretical justification, but is often used as a convenient characterization of an *a priori* unknown distribution of relaxation rates. As in SDRO we would expect different internal fields for muons stopping on the DyO₆ octahedral site and on the RuO₆ octahedral site. The present approach to fitting the data with the use of a stretched exponential function therefore seems appropriate. We have also fitted the μ SR spectra of SDRO using a simple exponential function (see the Supplemental Material, Figs. S1 and S2 [57]) and the results are very similar to that obtained from a stretched exponential function fit given here. It is to be noted that a simple exponential function fits better to the low temperature and high temperature data, however the fits are not very good for the data close to T_N .

Figure 7(b) shows that at high temperature (i.e., above the Néel temperature) the μ SR spectra exhibit a moderate relaxation rate, which is due to the spin fluctuation from the Ru⁵⁺ and Dy³⁺ moments. With decreasing temperature, the relaxation rate increases and exhibits a peak near 40 K, followed by a rapid loss of muon initial asymmetry [Fig. 7(a)] below 40 K. Between 1.2 and 40 K, the asymmetry loss is almost 2/3. In the polycrystalline sample, we expect three components of internal field at the muon stopping sites. The loss of 2/3 can be understood as implanted muons see three components of the internal field in the polycrystalline sample. One longitudinal component, which is the component of internal field that is parallel to the incident muon beam (1/3 component), will give a relaxation and no oscillations. While the remaining two components seen by the muons (2/3 component) are the transverse components (i.e., components of the internal field perpendicular to the incident muon beam), which give oscillations. As the internal field from the Dy³⁺ and Ru⁵⁺ moments are expected to be larger and inhomogeneous, the

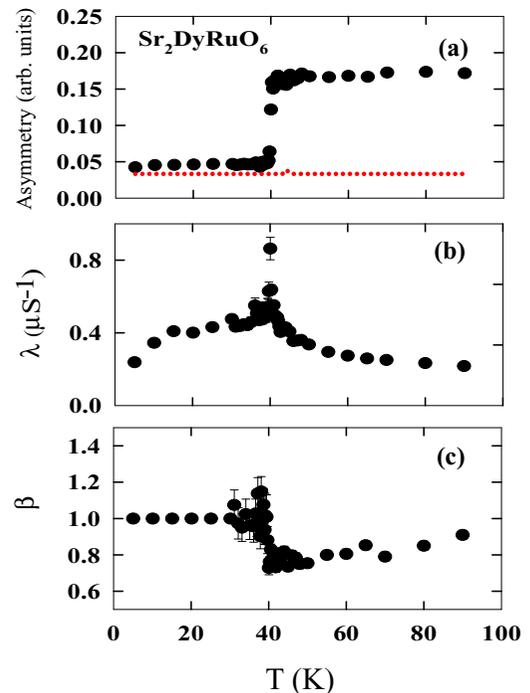


FIG. 7. The temperature dependent fit parameters obtained from the zero-field μ SR spectra. (a) Initial muon asymmetry versus temperature, please note that the symbols are the sample asymmetry only after subtracting the background asymmetry from the sliver sample holder shown by the dotted line, see Eq. (1). (b) Relaxation rate versus temperature and (c) exponent β versus temperature. The dotted line in (a) shows the temperature independent background from the sample holder.

oscillations in the muon time evolution asymmetry will be damped rapidly. Therefore, due to the pulse width of the ISIS muon beam, it is difficult to observe these oscillations in the 2/3 component as this signal damps very quickly in the short time window close to zero time. We therefore attribute the observed jump in A_0 at 40 K to long-range magnetic ordering of both Dy³⁺ and Ru⁵⁺ moments, as also observed in the heat capacity and magnetization measurements. This is also in agreement with our finding from the neutron diffraction study discussed below. Furthermore, if we look in detail at the behavior of $\lambda(T)$ near 36 K, then there is weak evidence of a second phase transition in $\lambda(T)$. However, as expected, A_0 does not reveal any sign of a second transition as the system is already in a complete long-range magnetic ordered state below 40 K and hence cannot lose further asymmetry. We therefore attribute the weak change in the $\lambda(T)$ near 36 K as due to a very small change in the spin configuration. The exponent β reveals a temperature dependent behavior. It starts to decrease with temperature from ~ 1 at 90 K to ~ 0.7 at 40 K and then exhibits a sharp jump to 1 below T_{N1} and remains the same to lower temperature. It is interesting to note that a very recent μ SR study on La₂NaRuO₆ and La₂NaOsO₆ [40] reports similar features in $\lambda(T)$, while only exhibiting a single magnetic transition. The μ SR spectra of these compounds were also better fitted to a stretched exponential function. In SDRO there are no clear signs of frequency oscillations. This indicates that the internal fields at the muon

stopping sites, which are most likely close to the oxygen ions due to the negative charge on O^{2-} and positive charge on muon, are larger than 800 G. The width limit of the ISIS muon pulse does not allow the collection of data at very short times close to zero. The ordered moment of Dy^{3+} is quite high compared to Ru^{5+} (discussed in the neutron diffraction section), and it is highly possible that muons stopping in the DyO_6 octahedra will see a higher internal field than those stopping in the RuO_6 octahedra. We have also performed LF measurements at 5, 35, and 50 K up to 2.5 kG field (see Fig. S3 of the Supplemental Material (SM) [57]). At 5 and 35 K the initial asymmetry gradually recovers to 0.12 and 0.15, respectively, at 2.5 kG field. Furthermore, the relaxation rate also exhibits a gradual decrease with applied field. On the other hand, at 50 K the initial asymmetry does not change much with applied field, but the relaxation rate decreases slowly with applied field.

D. Neutron powder diffraction studies

To explore the origin of the two anomalies observed in various experiments, temperature dependent NPD measurements have been performed on the high intensity diffractometer D20 between 1.7 and 50 K with $\lambda = 1.594 \text{ \AA}$ (Fig. 8). Long scans of 45 min have been performed at 1.7, 10, 20, 30, 40 and 50 K and shorter scans of ~ 11 min were recorded between these temperatures while raising the temperature with a constant ramp of 1 K/5 min. Figure 8(a) represents the NPD patterns at 1.7 and 50 K, along with the difference curve plotted in green. Several new peaks emerge and some existing ones increase in intensity at low temperature indicating the onset of long-range magnetic ordering. The magnetic reflections can be indexed with $h + k + l$ being odd. Figure 8(b) is the 3D plot of the temperature dependent D20 data in a limited 2θ range in which the magnetic peaks are indexed using the nuclear unit cell, i.e., the magnetic propagation vector $k = (0, 0, 0)$. This is consistent with *A*-type antiferromagnetic order where ferromagnetic sheets within the *ab* plane are antiferromagnetically stacked along the long *c* axis with both the magnetic and crystal unit cell of same size. The presence of the (001) reflection indicates that the magnetic moments have components in the *ab* plane. The strongest magnetic peak represents in fact two reflections which can be indexed as (010)/(100), but cannot be resolved due to the $a \sim b$ pseudosymmetry. Magnetic symmetry analysis for the two possible magnetic sites of Dy^{3+} on $2c$ and Ru^{5+} on $2d$ was done in space group $P2_1/n$ with $k = (0, 0, 0)$ using the program BASIREPS [58,59]. Both sites possess the same two allowed irreducible representations (IR) with each having three basis vectors (BV) (Table I). The two IRs are differing in the directions of the ferromagnetic (F) or antiferromagnetic (AF) couplings between the symmetry related sites within the unit cell. While IR1 allows an AF coupling along the unit cell *b* direction with F couplings allowed along *a* and *c* directions. The IR2 describes just the opposite with a F coupling along *b* and AF couplings along *a* and *c*. Testing both IRs, it becomes clear that due to the above mentioned pseudosymmetry, it is possible to refine the low temperature magnetic structure model assuming either an AF coupling along *a* (IR1) or along *b* (IR2). In both cases a single BV is sufficient to

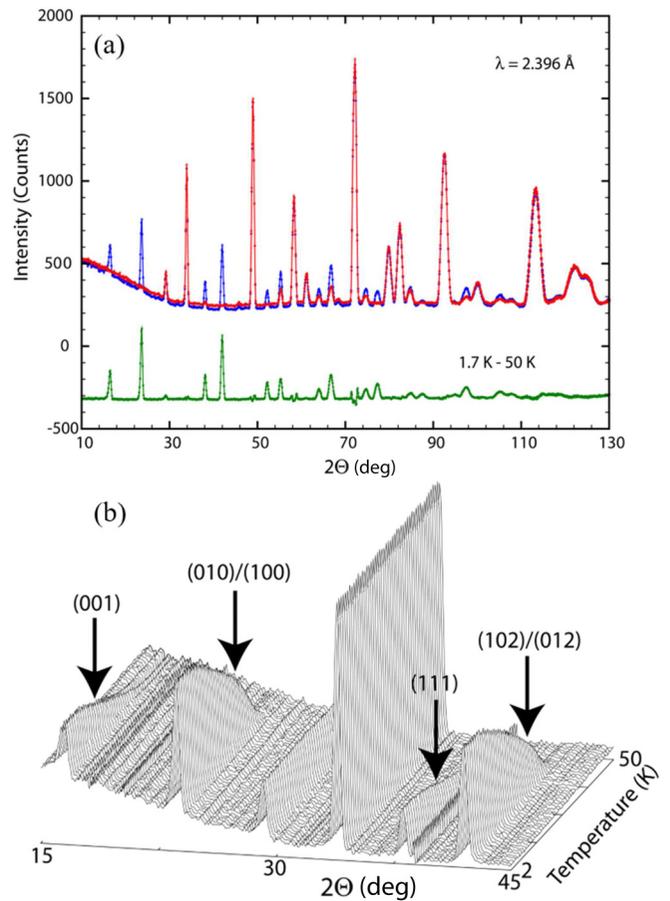


FIG. 8. Neutron diffraction pattern at 1.7 (blue curve) and 50 K (red curve) measured using D20. At 1.7 K, extra peaks are present compared to 50 K, which is due to magnetic ordering of the Dy^{3+} and Ru^{5+} moments. The green line at the bottom represents the difference curve (1.7–50 K) and shows only the magnetic Bragg peaks. (b) Thermal evolution of magnetic peak profiles between 2 and 50 K. The arrows show the magnetic Bragg peaks.

refine the magnetic structure, there is no clear indication of a ferromagnetic contribution to the magnetic scattering. As even our high resolution data from D2B at 3 K using $\lambda = 1.594 \text{ \AA}$ are not able to differentiate between the two models [Fig. 9(b)], additional NPD data have been collected at 3.5 K as well on the high-resolution powder diffractometer D2B using the longer wavelength of $\lambda = 2.399 \text{ \AA}$ in order to determine whether the moments are pointing in the *a* or *b* directions (Fig. 10). The only magnetic peak which allows us to differentiate between these two possibilities is the (100)/(010) doublet. Figure 10 shows the Rietveld refined NPD pattern measured at 3.5 K. The insert shows the enlarged view of the (010)/(100) peaks. It clearly shows that there is magnetic intensity at the (100) reflection but not on the (010) reflection, which confirms that the magnetic moments are pointing in the *b* direction. The resultant magnetic structure is shown in Fig. 11 in which the Ru^{5+} moments are represented in green color and the Dy^{3+} moments in red color. Using this model, the magnetic structure was refined using the high resolution data collected on D2B at 3 K and $\lambda = 1.595 \text{ \AA}$, with the magnetic intensity modeled as a second phase containing

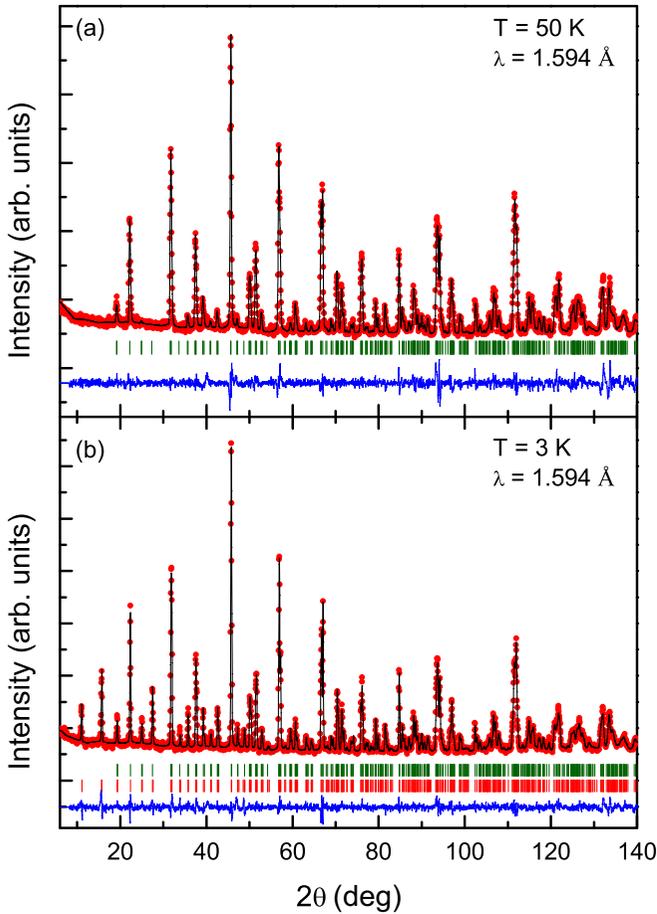


FIG. 9. Rietveld fit to NPD patterns collected at (a) 50 and (b) 3 K with $\lambda = 1.594 \text{ \AA}$ on D2B diffractometer at ILL: Red circles show the observed data points; the black line the calculated profile; the lower blue line the difference profile (obs. – calc.); upper green vertical markers indicate nuclear Bragg peak positions; lower red vertical markers in (b) indicate magnetic Bragg peak positions.

only the Dy and Ru atoms. The magnetic form factor used for Ru^{5+} is the one determined empirically in [60]. Figure 9(b) shows this refinement where the lower set of red tick marks correspond to the magnetic Bragg reflections. Table II shows the relevant bond lengths and angles at $T = 3 \text{ K}$ together with those determined at $T = 50 \text{ K}$ [Fig. 9(a)]: Cooling through the magnetic transition leads to a compression of the RuO_6 octahedra, whereas the DyO_6 octahedra elongate.

A cyclic structure refinement using the temperature dependent data from D20 allowed the determination of the thermal dependence of the magnetic moments of Dy^{3+} and Ru^{5+} and is shown in Fig. 12(a) with the normalized moments plotted in Fig. 12(b). It can be seen that for the present system, the Ru^{5+} moment attains saturation at a faster rate near $\sim 20 \text{ K}$ compared to Dy^{3+} which attains saturation only well below $\sim 10 \text{ K}$. This behavior is similar to other members of this family, like $\text{Sr}_2\text{HoRuO}_6$ and $\text{Sr}_2\text{TbRuO}_6$ [38]. However, the direction of the magnetic moments of Dy^{3+} and Ru^{5+} are different in the present system from those of the Ho and Tb based double ruthenates. While both the rare earth and the Ru^{5+} moments are along the c axis in $\text{Sr}_2\text{HoRuO}_6$, they are canted

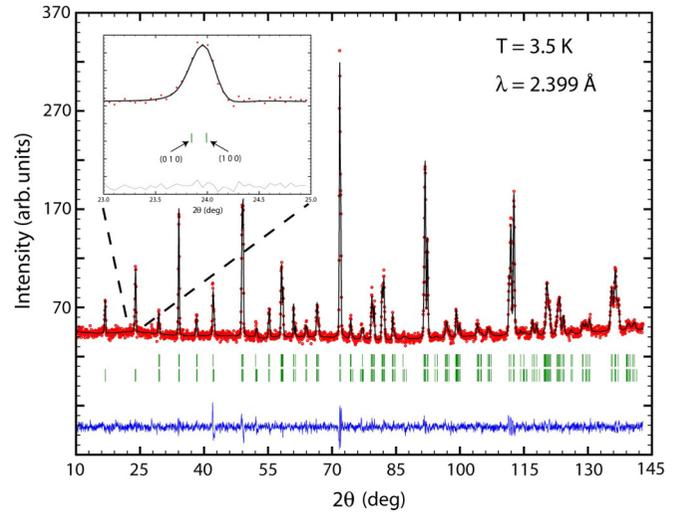


FIG. 10. Rietveld fit to NPD patterns collected at 3.5 K with $\lambda = 2.399 \text{ \AA}$ on D2B diffractometer at ILL: Red circles show the observed data points; the black line the calculated profile; the lower blue line the difference profile (obs. – calc.); upper green vertical markers indicate nuclear Bragg peak positions; lower red vertical markers indicate magnetic Bragg peak positions. The inset enlarges the view near 24° to highlight the magnetic peak (010)/(100).

by 20° from the c axis for $\text{Sr}_2\text{TbRuO}_6$ [38]. The moments of Ru and Pr in the ac plane (i.e., tilted away from the c axis) were also found in $\text{Ba}_2\text{PrRuO}_6$ [55]. Furthermore, the magnetic structure of $\text{Sr}_2\text{ErRuO}_6$ shows Ru^{5+} and Er^{3+} moments are mainly aligned along the c axis of the structure, forming an angle of $\sim 6^\circ$ with the c axis in the case of the Er^{3+} sublattice and $\sim 15^\circ$ for the Ru^{5+} moment [61]. In the present studied system, both the Dy^{3+} and Ru^{5+} moments are pointing along the b axis. The values of the Dy^{3+} and the Ru^{5+} moments at 1.7 K are $\mu_{\text{Dy}^{3+}} = 4.92(10) \mu_B$ and $\mu_{\text{Ru}^{5+}} = 1.94(7) \mu_B$. The strong reduction of $\mu_{\text{Dy}^{3+}}$ compared to the expected value

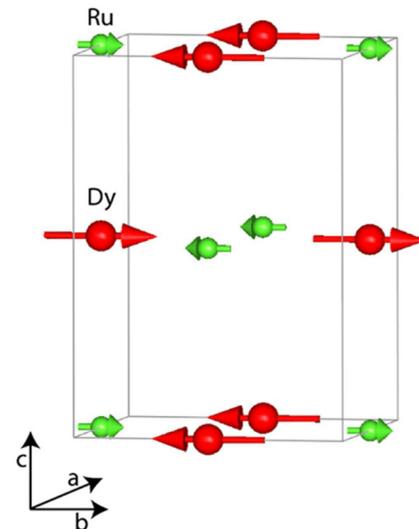


FIG. 11. The magnetic structure of $\text{Sr}_2\text{DyRuO}_6$ for $k = (0, 0, 0)$. The Dy^{3+} and Ru^{5+} moments are shown in red and green colors, respectively.

TABLE I. Basis vectors (BV) of the allowed irreducible representations (IR) for $k = (0, 0, 0)$ for the Wyckoff positions $2c$ (Dy) and $2e$ (Ru) of the space group $P2_1/n$.

IR1	BV1	BV2	BV3
x, y, z	100	010	001
$-x+\frac{1}{2}, y+\frac{1}{2}, -z+\frac{1}{2}$	-100	010	00-1
IR2			
x, y, z	100	010	001
$-x+\frac{1}{2}, y+\frac{1}{2}, -z+\frac{1}{2}$	100	0-10	001

of $10 \mu_B$ in line with similar discrepancies observed for the other rare earth Ru-based perovskites [38]. The Ru^{5+} moment value is similar to those reported for other members of this family [36,38,46] and points to the fact that in these systems the super-super exchange Ru-O-O-Ru interactions are the strongest magnetic interactions and control the Ru ordering. The similar values of $\mu_{\text{Ru}^{5+}}$ and of the magnetic transition temperatures in the different $\text{Sr}_2\text{LnRuO}_6$ systems [38] are explained by the weakness of the Ln-O-Ru interactions. As exemplified by the very low magnetic transition temperature of Dy_2O_3 ($T_N = 1.2$ K), Dy-O-Dy interactions are in general very weak. In the well-ordered double perovskite SDRO, only weaker super-super exchange Dy-O-O-Dy interactions are present which cannot be the origin of the primary ordering of Dy^{3+} at 39 K as supported by the order parameters given in Fig. 12(b). The temperature dependence of the Ru^{5+} moments exhibits a mean field power-law behavior with a critical exponent $\beta = 0.56(1)$, whereas that of the Dy^{3+} moments deviates from the power law. Hence, it appears that Ru^{5+} induces the rare earth ordering in these systems leading to a simultaneous ordering of Dy^{3+} at the same temperature, as also previously reported for $\text{Sr}_2\text{LnRuO}_6$ ($\text{Ln} = \text{Ho}$ and Tb) [38] as well as in R_2RuO_5 [62–64]. It should be noted that there appear neither appreciable changes in the magnetic peak profiles nor new magnetic Bragg peaks in the temperature dependent data across the second anomaly (~ 36 K). In particular we did not detect any additional superlattice peak in the temperature region between T_{N1} and T_{N2} , such as the one created by a propagation vector $k = (1/2, 1/2, 0)$ found by Bernardo *et al.*

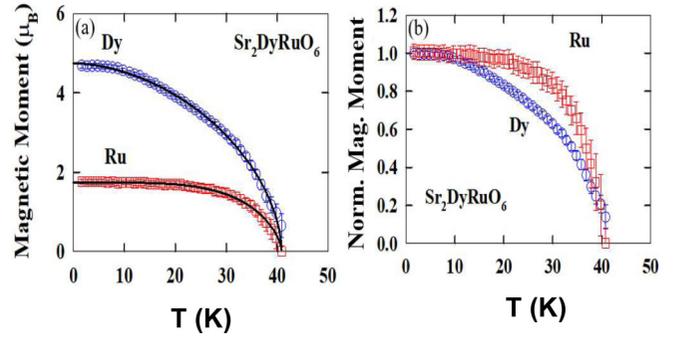


FIG. 12. Thermal variation of (a) Dy^{3+} and Ru^{5+} moments and (b) normalized moments of Dy^{3+} and Ru^{5+} .

[35] in Sr_2YRuO_6 . Furthermore, there is no indication for short-range correlations in the background below $T_{N1} = 40$ K. This indicates that the second transition might be associated with very small changes in the spin structure/spin reorientation which are beyond the detection limit even of the high intensity data. Although we are not able to detect any clear changes in the temperature dependent neutron data at T_{N2} , we tested the idea that a spin reorientation could be linked to the appearance of a symmetry allowed ferromagnetic component. Trying to include such a component in the refinement of the 1.7 K high intensity data, values of $0.8(5) \mu_B$ for the Dy site or $1.0(0.7) \mu_B$ for the Ru site were determined. Allowing a simultaneous presence of ferromagnetic components on both sites leads to the divergence of the refinements. Hence, and looking at the very large error bars, we can neither claim the presence of a ferromagnetic component nor its localization. Refinement of the long scan taken at 30 K including ferromagnetic components leads to even less precise values of about $0.4(0.8) \mu_B$ on either site. We therefore cannot ascertain but only mention that the second transition could be caused by the appearance of a ferromagnetic component leading to a spin reorientation. This ferromagnetic component would then be as well responsible for the weak ferromagnetic hysteresis observed in the magnetization isotherm at low temperature.

The possible magnetic frustration present in double perovskites is a consequence of the crystal structure as both B

TABLE II. Selected bond lengths (\AA) and bond angles (deg) across AFM ordering.

	50 K	3.5 K		50 K	3.5 K
Ru-O1	1.938(7)	1.952(6)	O1-Ru-O2	88.8(3)	89.8(3)
Ru-O2	1.948(7)	1.965(6)	O1-Ru-O3	89.3(3)	90.6(3)
Ru-O3	1.966(6)	1.954(5)	O2-Ru-O3	89.1(3)	89.7(3)
Dy-O1	2.232(7)	2.225(6)	O1-Dy-O2	92.1(3)	91.7(3)
Dy-O2	2.250(7)	2.238(6)	O1-Dy-O3	91.5(3)	89.3(2)
Dy-O3	2.217(6)	2.219(5)	O2-Dy-O3	91.3(3)	88.3(2)
Ru-O1-Dy	158.2(4)	157.3(4)	Ru-O3-Dy	155.5(4)	156.3(3)
Ru-O2-Dy	154.4(4)	155.0(4)	Ru-Dy	4.0883(1)	4.0832(1)
				4.0951(1)	4.0903(9)
Ru-Ru	5.7741(3)	5.7675(1)	Dy-Dy	5.7741(3)	5.7675(1)
	5.7747(3)	5.7915(1)		5.7984(1)	5.7915(1)
	5.8088(3)	5.8017(1)		5.8088(3)	5.8017(1)

sites form lattices of edge-sharing tetrahedra resembling the arrangement in a face centered cubic lattice. Nearest neighbor (NN) distances within one tetrahedron and to the neighboring tetrahedra are very similar so that any antiferromagnetic interaction must lead to a situation where conflicting demands on the orientations of equidistant neighboring spins arise. This is especially true when there is only one magnetic sublattice present, i.e., as in Sr_2YRuO_6 , which can lead to magnetic structures reflecting inherent frustration [36]. In SDRO two magnetic sublattices are present, one from Ru and another from Dy. If the intersublattice interactions Dy-O-Ru via superexchange are important, the frustration created through the intrasublattice interactions (i.e., between Ru-Ru and Dy-Dy) can be overcome and the system can order in a long-range magnetic state. The importance of the intersublattice interaction is in the case of SDRO reflected in the fact that both sublattices order at the same temperature (see Fig. 12) and with the same magnetic propagation vector. Intersublattice superexchange interactions of very similar distances and Ru-O-Dy angles are found along the c direction of the unit cell and within the a - b plane (Fig. 11). There are, however, four Ru-O-Dy NN pathways within the a - b plane while there are only two along the c direction. Assuming that these superexchange interactions are of antiferromagnetic type, the magnetic structure adopted in SDRO (Fig. 11) shows that the interactions within the a - b plane are dominant. Keeping in mind that the intrasublattice NN Ru-O-Ru superexchange interactions are still the strongest of all interactions—determining the temperature of the magnetic transition T_{N1} , see above—an antiferromagnetic alignment between Ru sites in neighboring a - b layers is established. This in turn forces a ferromagnetic alignment between Ru and Dy sites in neighboring a - b layers. The magnetic structure adopted sees therefore two antiferromagnetic and two ferromagnetic NN intrasublattice interactions for both the Ru and the Dy tetrahedra. As the monoclinic distortion is very small the intrasublattice interaction distances and angles are, however, very similar and one has to conclude that some frustration should still be present in the compound. This is neither visible through the presence of some static diffuse scattering in the neutron diffraction data nor in the frustration index $f = |\Theta_{\text{CW}}|/T_N \sim 0.7$ but could be reflected in the low values of the magnetic moment of $\text{Dy}^{3+} = 4.9(1) \mu_B$ and even $\text{Ru}^{5+} = 1.94(7) \mu_B$.

In the case of SDRO the two magnetic transitions observed in the heat capacity are not discernible in the neutron diffraction patterns as a change of the magnetic structure across the two transitions. This is not uncommon as there are several materials where heat capacity shows two (or three) peaks, but neutron diffraction data do not show clear visible changes between these transitions [65–67]. In the case of CeIrGe_3 the neutron diffraction measurements could only detect a very small change in the magnetic propagation vector from $k = (0, 0, 0.6667)$ to $(0, 0, 0.688)$ across the three transitions defined through three peaks in the heat capacity [65]. A second example concerns the multiferroic compound Mn_3TeO_6 exhibiting two peaks in the heat capacity at $T_{N1} = 23$ K and at $T^* = T_{N2} = 21$ K, temperature below which a multiferroic state emerges. The neutron diffraction study on Mn_3TeO_6 reveals only one magnetic transition at 24 K and the same incommensurately modulated magnetic spin structure with

the propagation vector $k = (0, 0, 0.43)$ is kept in the whole temperature range from 5 to 24 K [66]. A third example is YbRu_2Ge_2 [67] where the heat capacity shows two magnetic transitions, $T_{N1} = 6.5$ K and $T_{N2} = 5.5$ K. Neutron data revealed here only a tiny change in the magnetic propagation vector [from $\mathbf{k} = (0.352, 0, 0)$ to $(0.350, 0, 0)$] without any indication of this transition in the intensity of the magnetic reflections. These examples show that the presence of a peak in the heat capacity data does not necessarily lead to a detectable change in the neutron data and the first and third examples show that tiny changes in the configuration can induce peaks in the specific heat data. In the case of $\text{Sr}_2\text{DyRuO}_6$ treated here, we can associate the first peak to the appearance of the antiferromagnetic order. The second peak in the heat capacity of $\text{Sr}_2\text{DyRuO}_6$ could be related to the appearance of additional ferromagnetic interactions leading to the spin reorientation.

E. Inelastic neutron scattering studies

It is very important to understand the origin of the reduced magnetic moment and the strength of anisotropy in SDRO. We therefore have performed inelastic neutron scattering (INS) measurements on SDRO at various temperatures. Figures 13(a)–13(d) show the color contour maps of the scattering intensity, energy transfer vs momentum transfer (Q), at various temperatures between 7 and 45 K with a neutron incident energy $E_i = 15$ meV. At 7 K a strong band of excitation can be seen near 3.25 meV, and weak scattering intensity near 5.63, 6.8, and 8.9 meV, which are more clear when presented in 1D intensity versus energy plots between $Q = 0$ and 2 \AA^{-1} [Figs. 13(e) and 13(f)]. At 30 K, the 3.25 meV excitation softens and it seems scattering intensity emerges out from $Q = 1.1 \text{ \AA}^{-1}$, which is the magnetic Bragg peak with index $(0, 1, 0) / (1, 0, 0)$. Increasing the temperature to 37 K, the inelastic scattering broadens and transforms into diffuse scattering with an energy width of ~ 8 meV. At 45 and 55 K (above T_N), we have seen the presence of weak and broad diffuse scattering in the elastic cut (see Fig. S4 in the SM [57]), which suggests the presence of magnetic frustration/short-range correlations above T_N in SDRO. As we did not see a clear sign of diffuse scattering in the diffraction data on D20, it may suggest that the diffuse scattering has a dynamic nature. From the data at 5 K, we confirm the presence of a spin wave at 7 K with a spin gap of ~ 3.25 meV and zone boundary energy of 8.9 meV. It is to be noted that the observed scattering could be also partly interpreted as Zeeman splitting of the Dy low energy crystal field excitation (CEF) levels below T_N . As there are in total four magnetic atoms per unit cell (2 Dy and 2 Ru) one expects four spin wave branches, which is in agreement with the observed four excitations below 10 meV [Figs. 13(e) and 13(f)]. Next we compare the value of the spin gap (defined as the peak position in the energy cut, q integrated close to the AFM zone center for the powder samples) observed in the present system with those reported in other double perovskites systems [68–71]. Our INS on Sr_2YRuO_6 reveals a spin gap of 5 meV [68], which agrees with the reported results [72]. Spin gaps of 1.8(8) and 6(1) meV have been observed in $\text{La}_2\text{LiRuO}_6$ and $\text{La}_2\text{LiOsO}_6$, respectively [69], 2.57(4) meV in $\text{La}_2\text{MgIrO}_6$, 2.09(3) in $\text{La}_2\text{ZnIrO}_6$ [70],

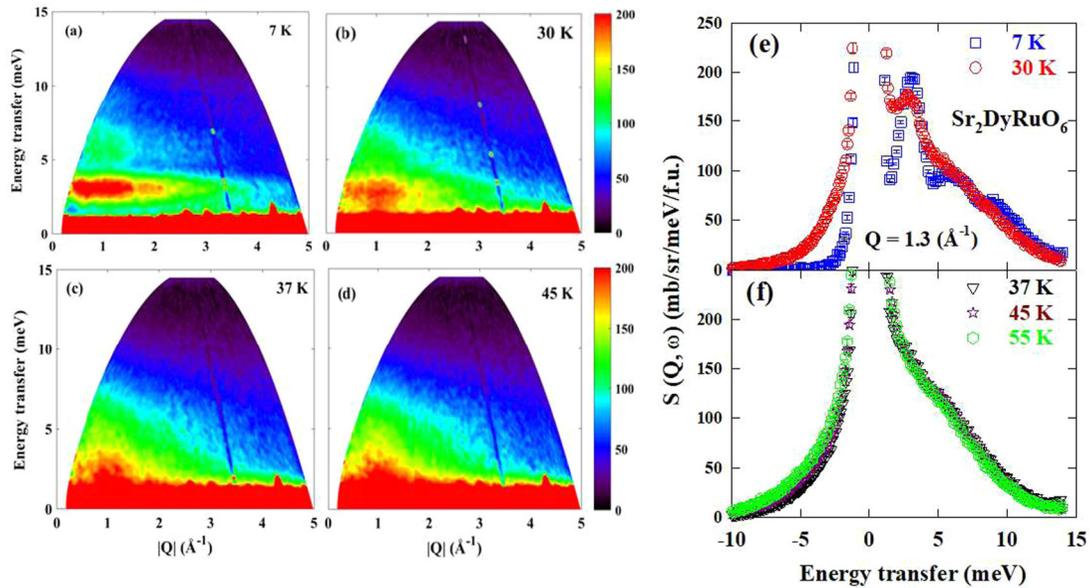


FIG. 13. (a)–(d) The color contour maps of scattering intensity versus momentum transfer at various temperatures measured with an incident energy of $E_i = 15 \text{ meV}$ on MERLIN. (e) and (f) The Q -integrated ($Q = 0$ to 2 \AA^{-1}) energy cuts at various temperatures between 7 and 55 K.

5 meV in Ba_2YRuO_6 [71], 2.75 meV in $\text{La}_2\text{NaRuO}_6$ [40], and 19 meV in $\text{Sr}_2\text{ScOsO}_6$ [73]. These results may suggest that the spin gap arises from the transition metal d electrons having strong spin-orbital coupling.

Now we discuss the crystal field excitations measured using $E_i = 130 \text{ meV}$ ($E_i = 250 \text{ meV}$ data are given in Fig. S6 of the SM). Figures 14(a) and 14(b) shows the color contour maps of the scattering intensity, energy transfer versus momentum transfer, at 7 and 100 K and Figs. 14(c)–14(e) shows the Q -integrated energy cuts from low- Q , medium- Q , and high- Q data. At 7 K and at low Q , strong intensity of scattering is observed near 46.6 and 73.8 meV. Furthermore,

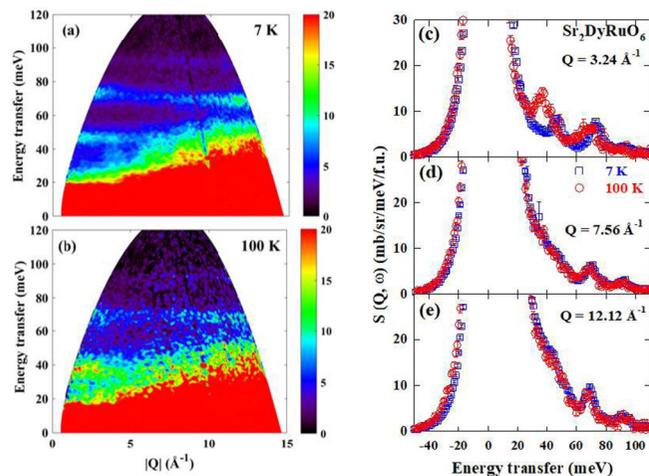


FIG. 14. (a) and (b) The color contour maps of scattering intensity versus momentum transfer at various temperature measured with an incident energy of $E_i = 130 \text{ meV}$ at 7 and 100 K on MERLIN. (c)–(f) The Q -integrated energy cuts at low Q , medium Q , and at high Q at 7 and 100 K.

a weak peak can be seen near 90.8 meV at higher Q , which has a lower intensity at lower Q . Plotting the Q dependence of the energy-integrated intensity of these peaks (see Fig. S5 in the SM [57]) the intensity of the peaks at 46.6 and 73.8 meV decreases initially and starts to increase at higher Q , while that of the 90.7 meV peak increases with Q . Moreover, the intensity of the 90.7 meV peak follows a Q^2 behavior [see the inset in Fig. S5(b)] as expected for phonon scattering. These observations indicate that the 46.6 and 73.8 meV peaks at low Q are due to the crystal field excitations of Dy^{3+} ion, while the 90.7 meV peak is due to purely phonon scattering. The increase of the intensity of the 46.3 and 73.8 meV peaks at higher Q indicates that phonon modes are present at similar energies suggesting the presence of phonon and CEF coupling. The assignment of CEF and phonon peaks seen in 130 meV data was also confirmed through the measurements with $E_i = 250 \text{ meV}$ at 7 and 120 K (see Fig. S6 of the SM [57]).

Looking at the data taken at 100 K with an incoming energy of 130 meV (250 meV at 120 K) it becomes clear that a new strong peak near 37.3 meV appears (same in 250 meV data at 120 K) and that the 73.8 meV peak has shifted to lower energy. It seems that the 46.6 meV peak has also shifted to lower energy. We attribute this new observed peak near 37.3 meV to the excited state transition from the CEF levels below 10 meV, which become populated at 100 K. The shift in the peaks could be due to various origins, i.e., magnetoelastic coupling and the Zeeman field at 7 K from the ordered state moments of the Dy and Ru.

We now discuss the CEF splitting of the Dy^{3+} ($4f^9$) ions based on the CEF Hamiltonian in order to provide further support for our interpretation of the INS spectra. The point symmetry of the Dy^{3+} ions is triclinic (1 or C_i) in the monoclinic $P2_1/c$ crystal structure of SDRO. In such a low symmetry, the CEF Hamiltonian requires 15 CEF parameters

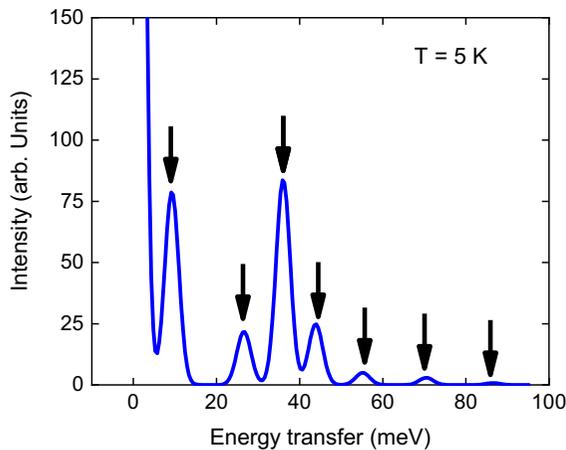


FIG. 15. Crystal field excitations calculated using a point charge model with positive sign of B_2^0 in Mantid-crystal field program [75] at 7 K. The linewidth of the excitations was taken as 4 meV. Seven crystal field excitations (doubly degenerate) can be seen. The overall splitting qualitatively agrees with the experimental data.

to be estimated from the INS spectra which is a difficult task. Considering an odd number of electrons, $4f^9$ of Dy^{3+} ions and Kramers' theorem [54], which says that for an odd number of electrons, the minimum degeneracy of CEF levels should be twofold (or doubly degenerate) in the paramagnetic state. We therefore expect that the $J = 15/2$ ground state (16-fold degeneracy = $2J + 1$) should split into 8 CEF levels with twofold degeneracy of each level above T_N . Furthermore, below the magnetic ordering these 8 doublets will split into 16 singlets. Hence if all CEF excitations from the ground state are allowed then in the paramagnetic state, we should expect a minimum of 7 CEF excitations from the ground state. In addition we will have additional CEF excitations due to the excited state transitions at $T = 100 \text{ K} > T_N$, if there exists low energy CEF levels and the excited state transitions are allowed. Given that we have observed only two CEF transitions at 46.6 and 73.8 meV (in addition to four spin wave type excitations below 9 meV) at $T = 7 \text{ K}$ and one additional excited state transition at 37.3 meV at 100 K a quantitative analysis of the INS data based on CEF model is not feasible. We have therefore used a point charge model to estimate the 15 CEF parameters (including 12 complex parameters this gives a total of 27 CEF parameters, see Table I in the SM [6]) of the CEF Hamiltonian and calculated INS spectra (Fig. 15) based on these estimated parameters. The simulated INS CEF spectrum at 7 K is given in Fig. 15. It is to be noted that the point charge model gave a negative sign of B_2^0 , corresponding to the c axis as the easy axis of the magnetization. This is in contradiction to the b -axis moment direction observed in the neutron diffraction data (Fig. 11). In order to account for this we have used a positive sign of B_2^0 , which does not change the calculated CEF excitations given in Fig. 15, but puts the moment in the ab plane. The simulation shows a qualitative agreement with the experimental data giving an overall CEF splitting of 86.6 meV, which is in good agreement with the observed CEF splitting of 73.8 meV. We have calculated the single crystal (along all three directions) susceptibility and powder averaged susceptibility (see Fig. S7

of the SM [57]) as well as moments of the Dy ion using the CEF parameters based on the point charge model at 2 K in an applied field of 7 T. The values of the Dy moment are $\mu_x = 7.5 \mu_B$, $\mu_y = 6.6 \mu_B$, $\mu_z = 0.85 \mu_B$, and the powder averaged moment $\mu_{\text{poly}} = 4.98 \mu_B$. The calculated powder average moment value of the Dy ion is slightly larger than the value of $4.25 \mu_B$ (total moment from the Dy and Ru) at 2 K and 7 T observed in the experimental data, which shows the limitation of the simple point charge model which does not include any covalence effects.

IV. CONCLUSIONS

Our combined heat capacity, magnetization, μSR , neutron diffraction, and inelastic neutron scattering results demonstrate that $\text{Sr}_2\text{DyRuO}_6$ (SDRO) exhibits a long-range ordered magnetic ground state below 40 K. The heat capacity reveals a clear sign of two magnetic transitions, which are also indirectly supported through the magnetic susceptibility (both AC and DC) measurements. Our μSR and neutron diffraction studies provide direct support for long-range magnetic ordering below 40 K. The neutron diffraction study shows that all the observed magnetic Bragg peaks between 2 and 40 K can be indexed using the magnetic propagation vector $k = (0, 0, 0)$. The magnetic structure shows that both the Dy and Ru atoms are arranged in type-I antiferromagnetic structure, which consists of interpenetrating sublattices of Dy^{3+} and Ru^{5+} atoms. In the ab plane, the Dy^{3+} and Ru^{5+} moments are aligned AFM to each other, while along the c axis they show FM coupling. Interestingly, the magnetic ordering is primarily governed by the $4d$ moment on the Ru^{5+} atoms and the Dy^{3+} moments follow the Ru ordering at the same temperature (T_{N1}). In addition, it appears that the interactions responsible for the Dy ordering are weaker than the interactions responsible for the Ru ordering. The origin of the second anomaly in the heat capacity still remains an open question as the neutron diffraction study shows only one magnetic transition at 40 K and further single crystal neutron diffraction study will be important to understand the origin of the two phase transitions in SDRO. A possible cause for this second transition could reside in the appearance of an additional ferromagnetic component at T_{N2} . From the inelastic neutron scattering study, we have estimated the spin gap of 3.25 meV in the spin wave spectrum with a maximum zone boundary energy of 8.9 meV. Furthermore, we have also discussed the presence of crystal field excitations and their role in the observed reduced moment of the Dy^{3+} ions estimated through the neutron diffraction. The total CEF splitting observed in the experimental data agrees very well with that calculated using the point charge model for the Dy^{3+} ion. The present work will generate interest in the condensed matter theory community to develop a realistic model to find out the common origin of the two magnetic phase transitions in the double perovskite family.

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